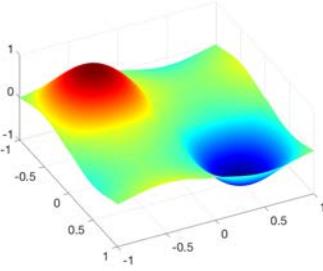


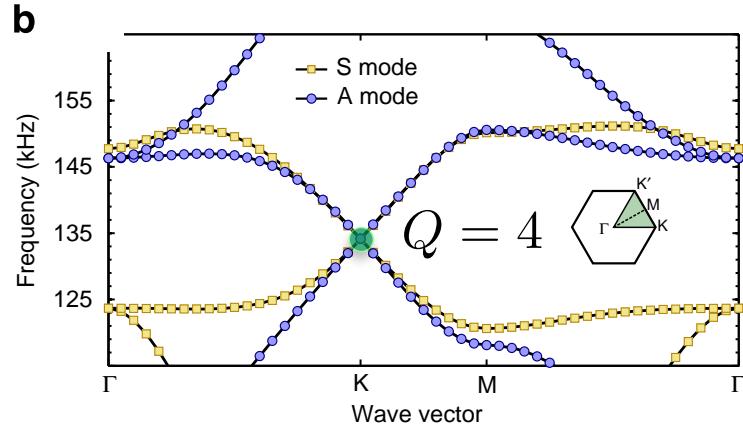
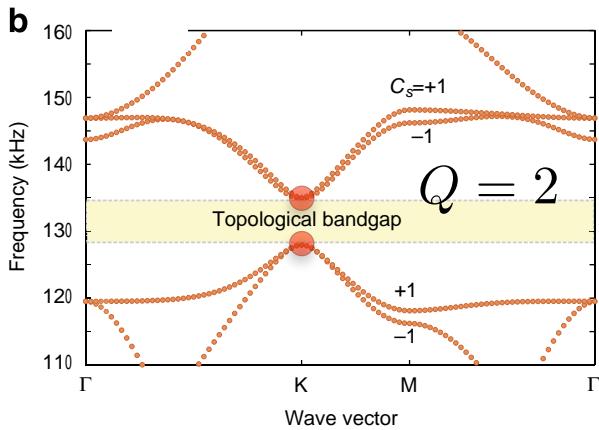
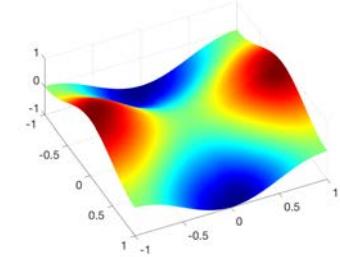
Part II

Dirac, Dirac-like, and almost-Dirac points

Repeated eigenvalues



Q projections (phonons)



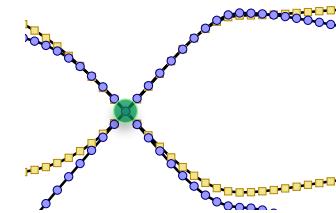
$$O(\epsilon^{-2}): \quad -\tilde{\lambda}_n^{\mathbf{a}} \rho \tilde{w}_0 - \nabla \cdot (G \nabla \tilde{w}_0) = 0 \quad \text{in } Y_{\mathbf{a}}$$

$$\tilde{w}_0(\mathbf{x}) = w_0 \tilde{\varphi}_n^{\mathbf{a}}(\mathbf{x})$$

$$\sum_q \underline{w}_{0q} \tilde{\varphi}_{nq}^{\mathbf{a}}(\mathbf{x})$$

Mousavi et al. (2015)

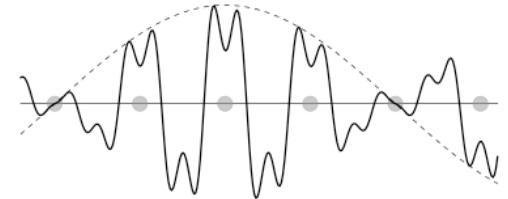
Energy



$$(\rho \tilde{\varphi}_{np}^{\boldsymbol{a}}, \tilde{\varphi}_{nq}^{\boldsymbol{a}})_{Y_{\boldsymbol{a}}} = 0 \quad \text{for } p \neq q$$

Bloch wave

$$u(\boldsymbol{x}) = \tilde{u}(\boldsymbol{x}) e^{i\boldsymbol{k}\cdot\boldsymbol{x}}, \quad \tilde{u} : Y_{\boldsymbol{a}}\text{-periodic}$$



Mean energy density Willis (2016) J MPS, 97

$$\bar{E} = E_p + E_k = \frac{1}{2} \omega^2 (\rho \tilde{u}, \tilde{u})_{Y_{\boldsymbol{a}}}$$

Repeated eigenvalue

$$\tilde{w}_0(\boldsymbol{x}) = \sum_q \underline{w}_{0q} \tilde{\varphi}_{nq}^{\boldsymbol{a}}(\boldsymbol{x}) \quad \Rightarrow$$

Quanta of energy

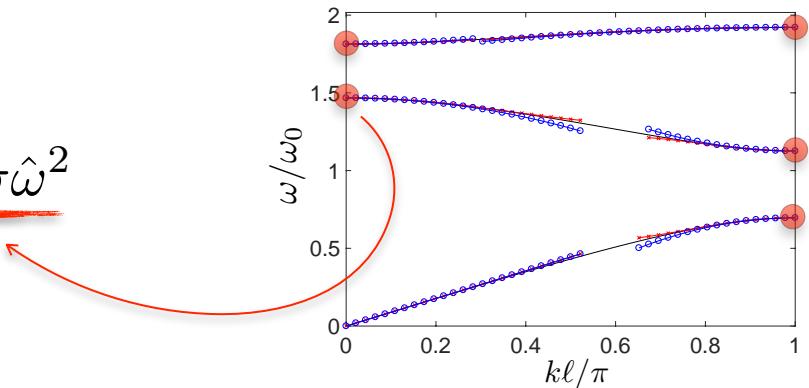
$$\rho_p^{(0)} = \langle \rho \tilde{\varphi}_{np}^{\boldsymbol{a}} \rangle_{\boldsymbol{a}}^{p\varphi}$$

$$\bar{E}_0 = \frac{1}{2} \omega^2 (\rho \tilde{u}_0, \tilde{u}_0) = \frac{1}{2} \omega^2 \tilde{f}^2 \sum_q \underline{\rho}_q^{(0)} w_{0q}^2$$

Scaling

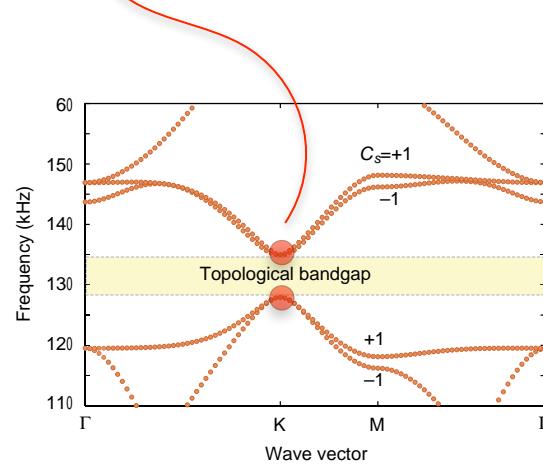
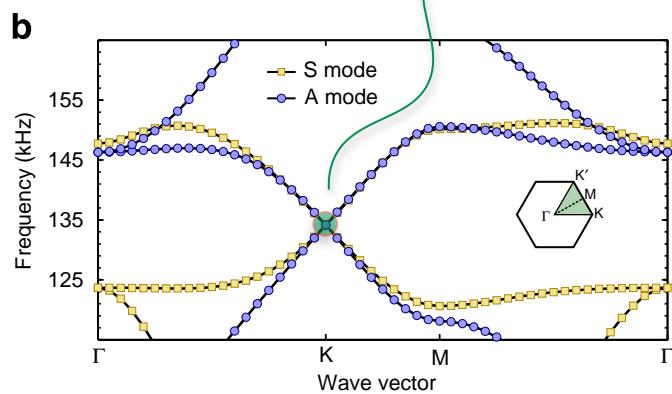
Isolated eigenvalues

$$\mathbf{k} = \mathbf{k}^a + \epsilon \hat{\mathbf{k}}, \quad \omega^2 = \tilde{\lambda}_n^a + \underline{\epsilon^2 \sigma \hat{\omega}^2}$$



Repeated eigenvalues

$$\mathbf{k} = \mathbf{k}^a + \epsilon \hat{\mathbf{k}}, \quad \omega^2 = \tilde{\lambda}_n^a + \underline{\epsilon \sigma \check{\omega}^2} + \underline{\epsilon^2 \sigma \hat{\omega}^2}$$



Cone or no cone?

Ansatz

$$O(\epsilon^{-2}): -\tilde{\lambda}_n^{\mathbf{a}} \rho \tilde{w}_0 - \nabla \cdot (G \nabla \tilde{w}_0) = 0 \quad \text{in } Y_{\mathbf{a}} \quad \Leftrightarrow \quad \tilde{w}_0(\mathbf{x}) = w_0 \tilde{\varphi}_n^{\mathbf{a}}(\mathbf{x})$$

$O(\epsilon^{-1})$: \rightarrow Identity

$O(\epsilon^0)$: \rightarrow Effective equation

$$\sum_q w_{0q} \tilde{\varphi}_{nq}^{\mathbf{a}}(\mathbf{x})$$

$$\sum_q \theta_{pq}^{(0)} \cdot (i\hat{\mathbf{k}}) w_{0q} + \underline{\sigma \check{\omega}^2} \rho_p^{(0)} w_{0p} = 0, \quad p = \overline{1, Q}$$

$$\theta_{pq}^{(0)} = \langle G \nabla \tilde{\varphi}_{nq}^{\mathbf{a}} \rangle_{\mathbf{a}}^{p\varphi} - \langle G \nabla \tilde{\varphi}_{np}^{\mathbf{a}} \rangle_{\mathbf{a}}^{q\varphi}$$

Generalized EVP

imaginary antisymmetric

$$\sum_q A_{pq} w_{0q} - \underline{\tau} \sum_q D_{pq} w_{0q} = 0, \quad p = \overline{1, Q}$$

Effective models

$$\tau = -\sigma \check{\omega}^2$$

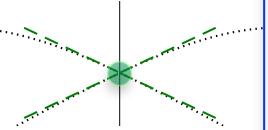
Generalized EVP

$$\sum_q A_{pq} w_{0q} - \tau \sum_q D_{pq} w_{0q} = 0, \quad p = \overline{1, Q}$$

$$A_{pq} = \theta_{pq}^{(0)} \cdot i\hat{\mathbf{k}}, \quad \theta_{pq}^{(0)} = \langle G \nabla \tilde{\varphi}_{nq}^{\mathbf{a}} \rangle_{\mathbf{a}}^{p\varphi} - \langle G \nabla \tilde{\varphi}_{np}^{\mathbf{a}} \rangle_{\mathbf{a}}^{q\varphi}$$

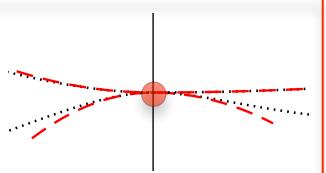
Full-rank A_{pq} : $w_{0q}=0$ & $\tau \neq 0$, system of 1st-order PDE's for w_{1q}

$$-\sum_q A_{pq} w_{1q} - \sigma \check{\omega}^2 \sum_q D_{pq} w_{1q} = \langle e^{i\mathbf{k}^{\mathbf{a}} \cdot \mathbf{x}} \rangle_{\mathbf{a}}^{p\varphi}, \quad p = \overline{1, Q}$$

Dirac 

Trivial A_{pq} : $w_{0q} \neq 0$ & $\tau = 0$, system of 2nd-order PDE's for w_{0q}

$$-\sum_q B_{pq} w_{0q} - \sigma \hat{\omega}^2 \sum_q D_{pq} w_{0q} = \langle e^{i\mathbf{k}^{\mathbf{a}} \cdot \mathbf{x}} \rangle_{\mathbf{a}}^{p\varphi}, \quad p = \overline{1, Q}$$

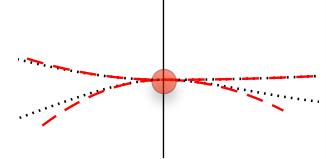
Wave 

$$B_{pq} = \mu_{pq}^{(0)} : (i\hat{\mathbf{k}})^2, \quad \mu_{pq}^{(0)} = \langle G \{ \nabla \chi_q^{(1)} + \mathbf{I} \tilde{\varphi}_{nq}^{\mathbf{a}} \} \rangle_{\mathbf{a}}^{p\varphi} - (G \{ \chi_q^{(1)} \otimes \nabla \tilde{\varphi}_{np}^{\mathbf{a}} \}, 1)_{\overline{Y}_{\mathbf{a}}}$$

Caution

Trivial A_{pq} : $w_{0q} \neq 0$ & $\tau = 0$, system of 2nd-order PDE's for w_{0q}

$$-\sum_q B_{pq} w_{0q} - \sigma \hat{\omega}^2 \sum_q D_{pq} w_{0q} = \langle e^{i\mathbf{k}^a \cdot \mathbf{x}} \rangle_a^{p\varphi}, \quad p = \overline{1, Q}$$



$$B_{pq} = \mu_{pq}^{(0)} : (i\hat{\mathbf{k}})^2, \quad \mu_{pq}^{(0)} = \langle G\{\nabla \chi_q^{(1)} + \mathbf{I}\tilde{\varphi}_{nq}^a\} \rangle_a^{p\varphi} - (G\{\chi_q^{(1)} \otimes \nabla \tilde{\varphi}_{np}^a\}, 1)_{\overline{Y}_a}$$

$$\tilde{w}_0(\mathbf{x}) = \sum_q w_{0q} \tilde{\varphi}_{nq}^a(\mathbf{x})$$

$$\chi_q^{(1)} \in (H_{p0}^{1a}(Y_a))^d$$

$$\tilde{w}_1(\mathbf{x}) = \sum_q w_{0q} \chi_q^{(1)}(\mathbf{x}) \cdot i\hat{\mathbf{k}} + \sum_q w_{1q} \tilde{\varphi}_{nq}^a(\mathbf{x})$$

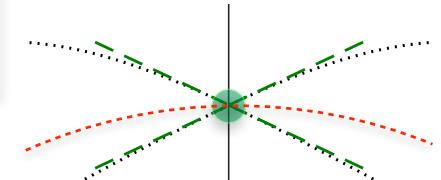
$$\tilde{\lambda}_n^a \rho \chi_q^{(1)} + \nabla \cdot (G(\nabla \chi_q^{(1)} + \mathbf{I}\tilde{\varphi}_{nq}^a)) + G \nabla \tilde{\varphi}_{nq}^a - \sum_r \frac{1}{\rho_r^{(0)}} \theta_{rq}^{(0)} \rho \tilde{\varphi}_{nr}^a = 0 \quad \text{in } Y_a, \quad q = \overline{1, Q}$$

$$\boldsymbol{\nu} \cdot G(\nabla \chi_q^{(1)} + \mathbf{I}\tilde{\varphi}_{nq}^a)|_{x_j=0} = -\boldsymbol{\nu} \cdot G(\nabla \chi_q^{(1)} + \mathbf{I}\tilde{\varphi}_{nq}^a)|_{x_j=(1+a_j)\ell_j}.$$

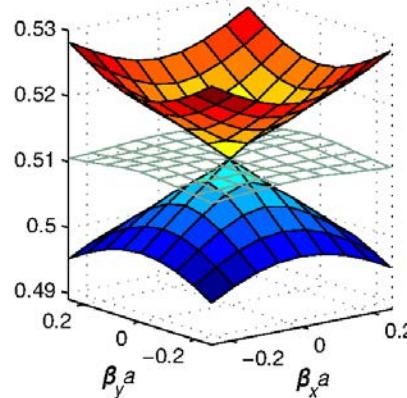
Mixed degeneracy

Partial rank A_{pq} ($Q-1$):

$$-(\mu_{11}^{(0)} : (i\hat{\mathbf{k}})^2 + \sigma\hat{\omega}^2\rho_1^{(0)})w_0 = \dots \quad 1 \text{ parabola}$$

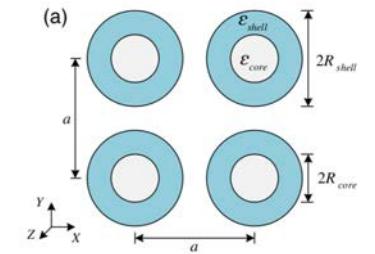


$$-\sum_{q=2}^Q A_{pq} w_{1q} - \sigma\check{\omega}^2 \sum_{q=2}^Q D_{pq} w_{1q} = \dots \quad Q-1 \text{ cones}$$



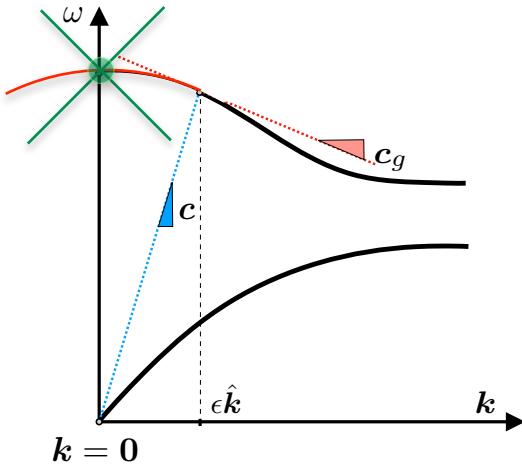
“Dirac-like cones”

Ashraf & Faryad (2015)
J. Nanophotonics, 9



Zero-index metamaterials

EM waves:
 $\epsilon=0$ and/or $\mu=0$



Phase velocity $c \rightarrow \infty$ (zero refraction index)

Applications to cloaking and shielding

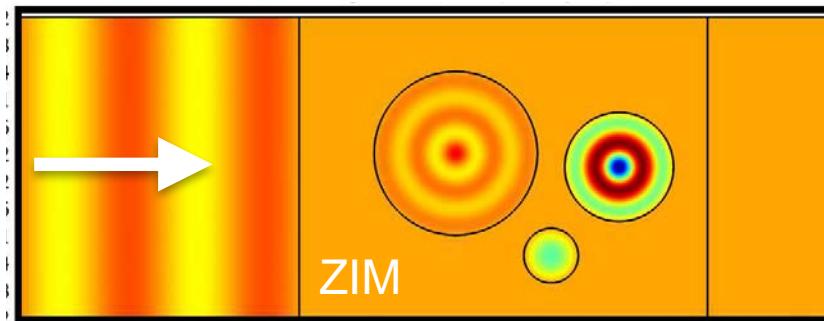
Ziolkowski (2004) *Phys. Rev. E*, **70**

Nguyen, Chen & Halterman (2010), *Phys. Rev. Lett.*, **105**

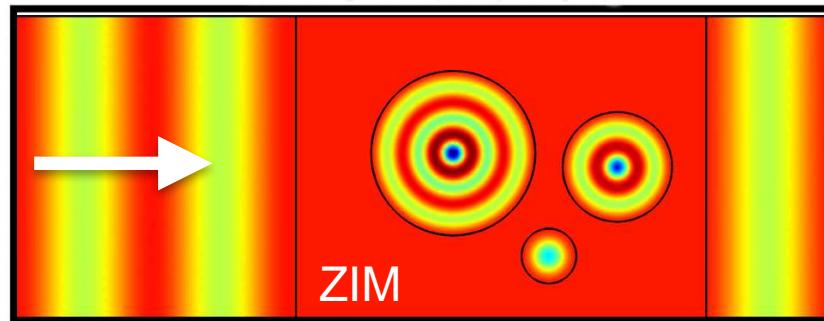
Chan, Hang & Huang (2012), *Adv. Optoelectronics*

Ashraf & Faryad (2015) *J. Nanophotonics*, **9**

Shielding

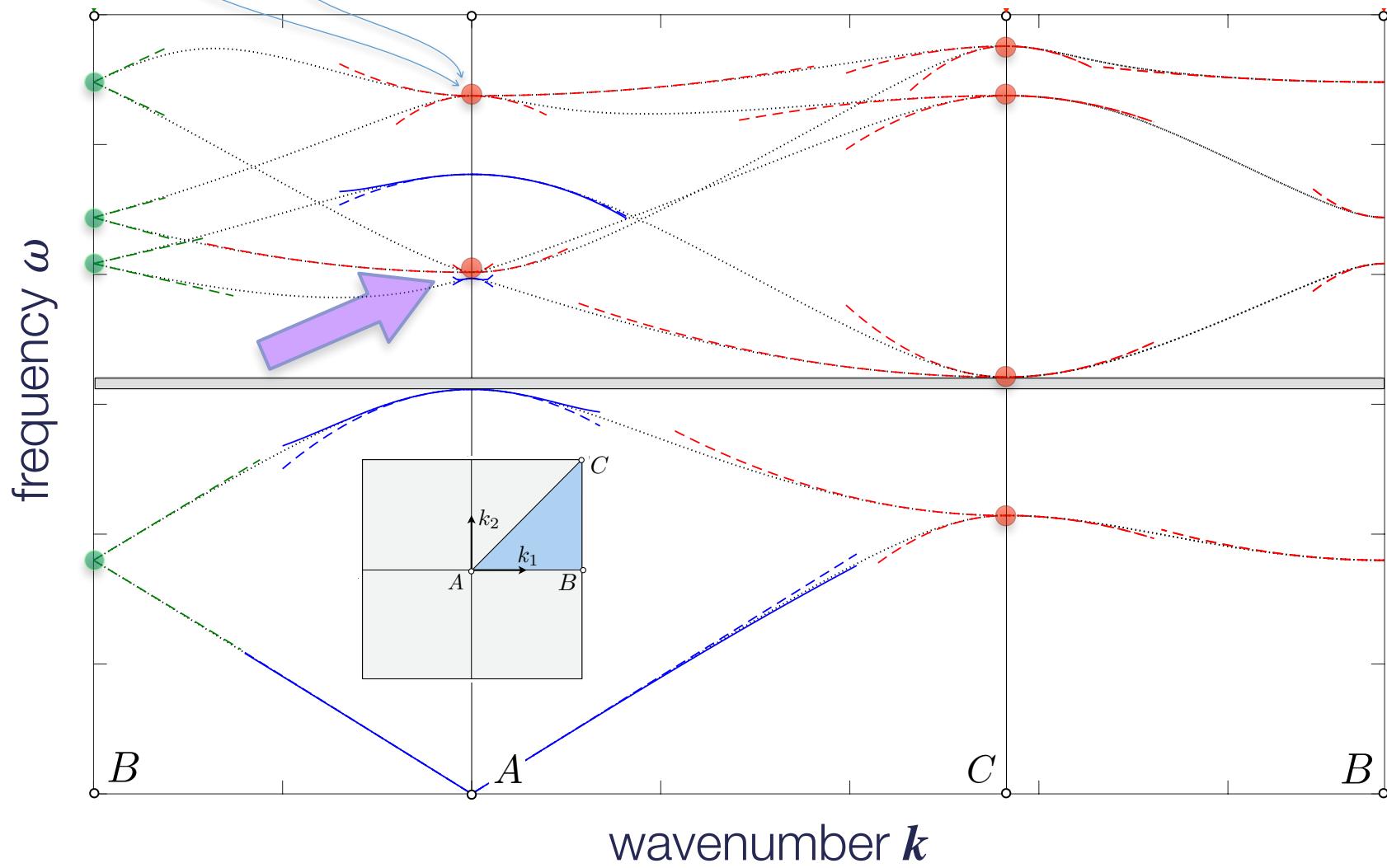
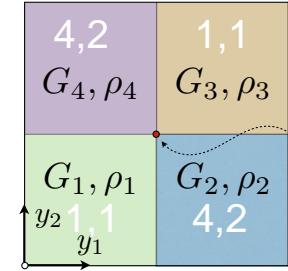
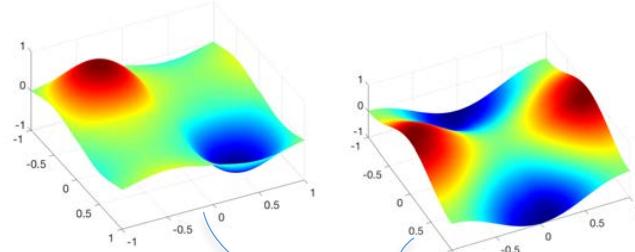


Cloaking



Nguyen, Chen &
Halterman (2010)

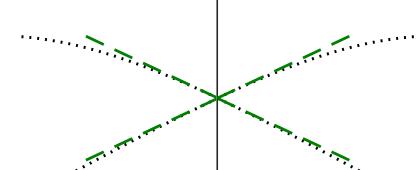
Chessboard



Clusters of nearby eigenvalues

Repeated eigenvalue

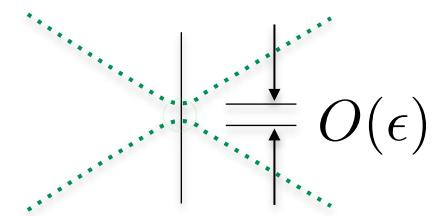
$$-\sum_q A_{pq} w_{1q} - \sigma \check{\omega}^2 \sum_q D_{pq} w_{1q} = \langle e^{i\mathbf{k}^a \cdot \mathbf{x}} \rangle_{\mathbf{a}}^{p\varphi}$$



Nearby eigenvalues?

$$\tilde{\lambda}_{j(q)}^{\mathbf{a}} = \tilde{\lambda}_n^{\mathbf{a}} - \underline{\epsilon \gamma_q}, \quad q = \overline{1, Q}$$

$$A_{pq} \rightarrow A_{pq}^{\gamma} = \theta_{pq}^{(0)} \cdot i\hat{\mathbf{k}} + \underline{\epsilon \gamma_q D_{pq}}$$



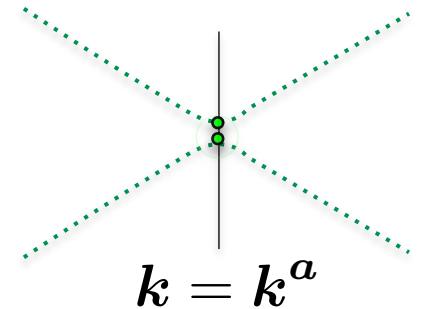
$$\begin{bmatrix} 0 & A_{12} \\ -A_{12} & 0 \end{bmatrix} \rightarrow \begin{bmatrix} 0 & A_{12} \\ -A_{12} & \gamma_2 \rho_2^{(0)} \end{bmatrix}$$

“blunted cones”

Cluster

Nearby eigenvalues

$$\begin{aligned}
 -\nabla \cdot (G \nabla \tilde{\varphi}_{nq}^{\boldsymbol{a}}) &= \tilde{\lambda}_n^{\boldsymbol{a}} \rho \tilde{\varphi}_{nq}^{\boldsymbol{a}} + \underline{\epsilon \gamma_q \rho \tilde{\varphi}_{nq}^{\boldsymbol{a}}} \quad \text{in } Y_{\boldsymbol{a}}, \\
 a_j \tilde{\varphi}_{nq}^{\boldsymbol{a}}|_{x_j=0} &= -a_j \tilde{\varphi}_{nq}^{\boldsymbol{a}}|_{x_j=\ell_j}, \\
 \boldsymbol{\nu} \cdot G \nabla \tilde{\varphi}_{nq}^{\boldsymbol{a}}|_{x_j=0} &= -\boldsymbol{\nu} \cdot G \nabla \tilde{\varphi}_{nq}^{\boldsymbol{a}}|_{x_j=(1+a_j)\ell_j}
 \end{aligned}$$



Governing equation

$$O(\epsilon^{-2}): \quad -\tilde{\lambda}_n^{\boldsymbol{a}} \rho \tilde{w}_0 - \nabla \cdot (G \nabla \tilde{w}_0) = 0 \quad \text{in } Y_{\boldsymbol{a}}$$

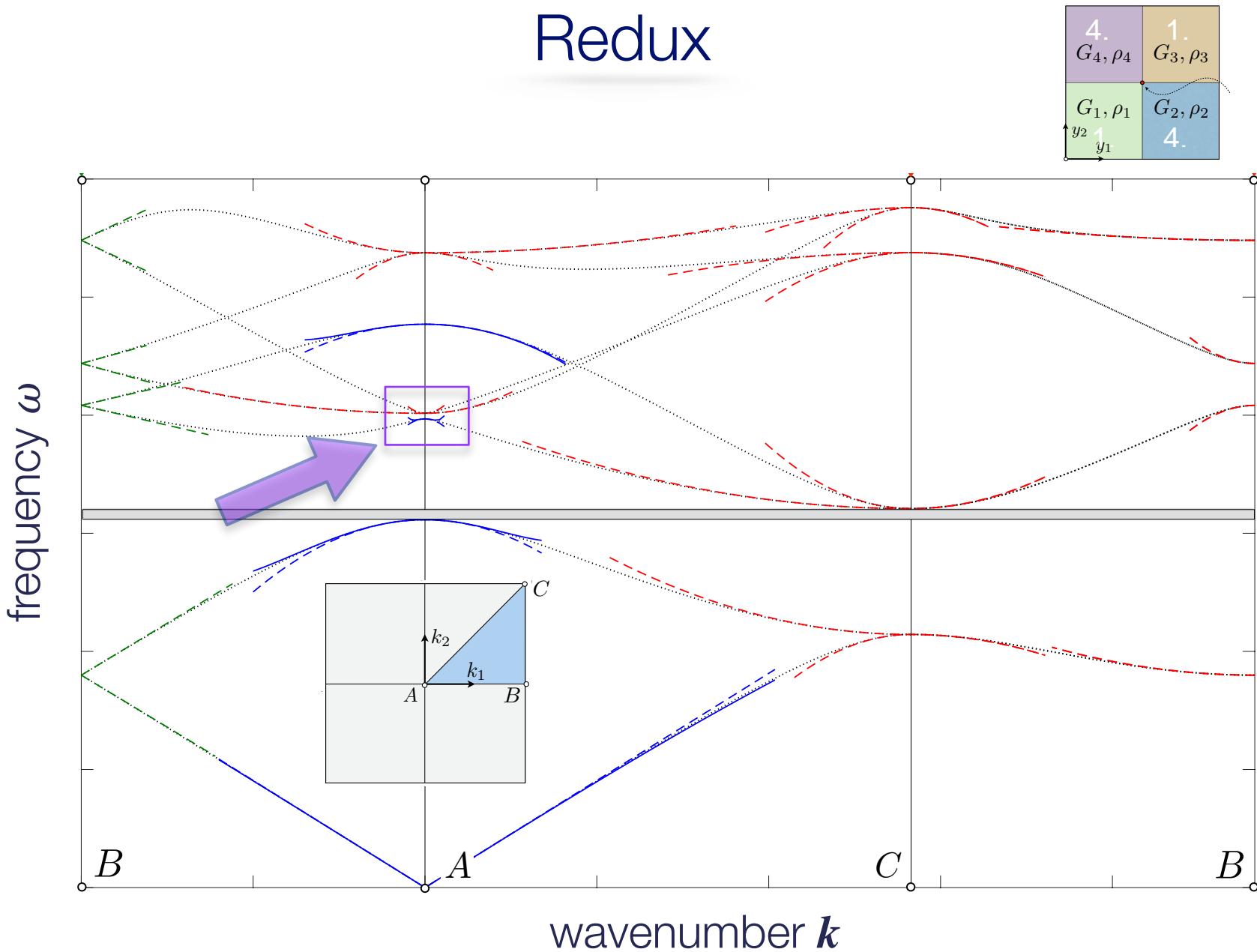
Leading-order solution

$$\tilde{w}_0(\boldsymbol{x}) = \sum_q w_{0q} \tilde{\varphi}_{nq}^{\boldsymbol{a}}(\boldsymbol{x})$$

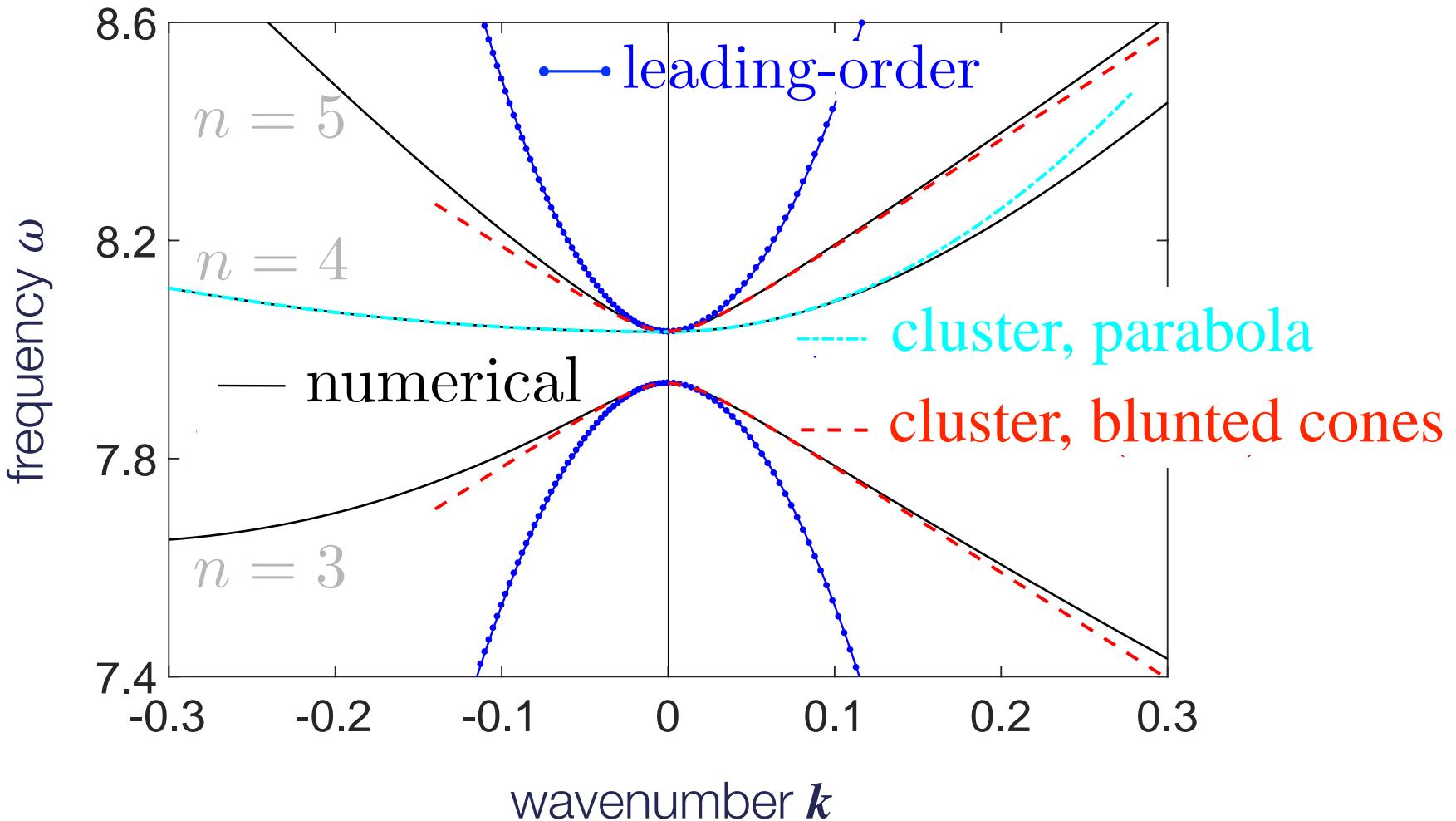
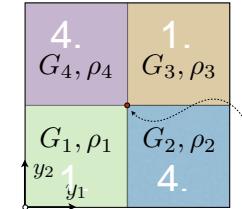
Residual

$$\epsilon \rho(\boldsymbol{x}) \sum_q \gamma_q w_{0q} \tilde{\varphi}_{nq}^{\boldsymbol{a}}(\boldsymbol{x}) \quad \xrightarrow{\hspace{1cm}} \quad A_{pq}^{\gamma} = \theta_{pq}^{(0)} \cdot i \hat{\boldsymbol{k}} + \underline{\gamma_q D_{pq}}$$

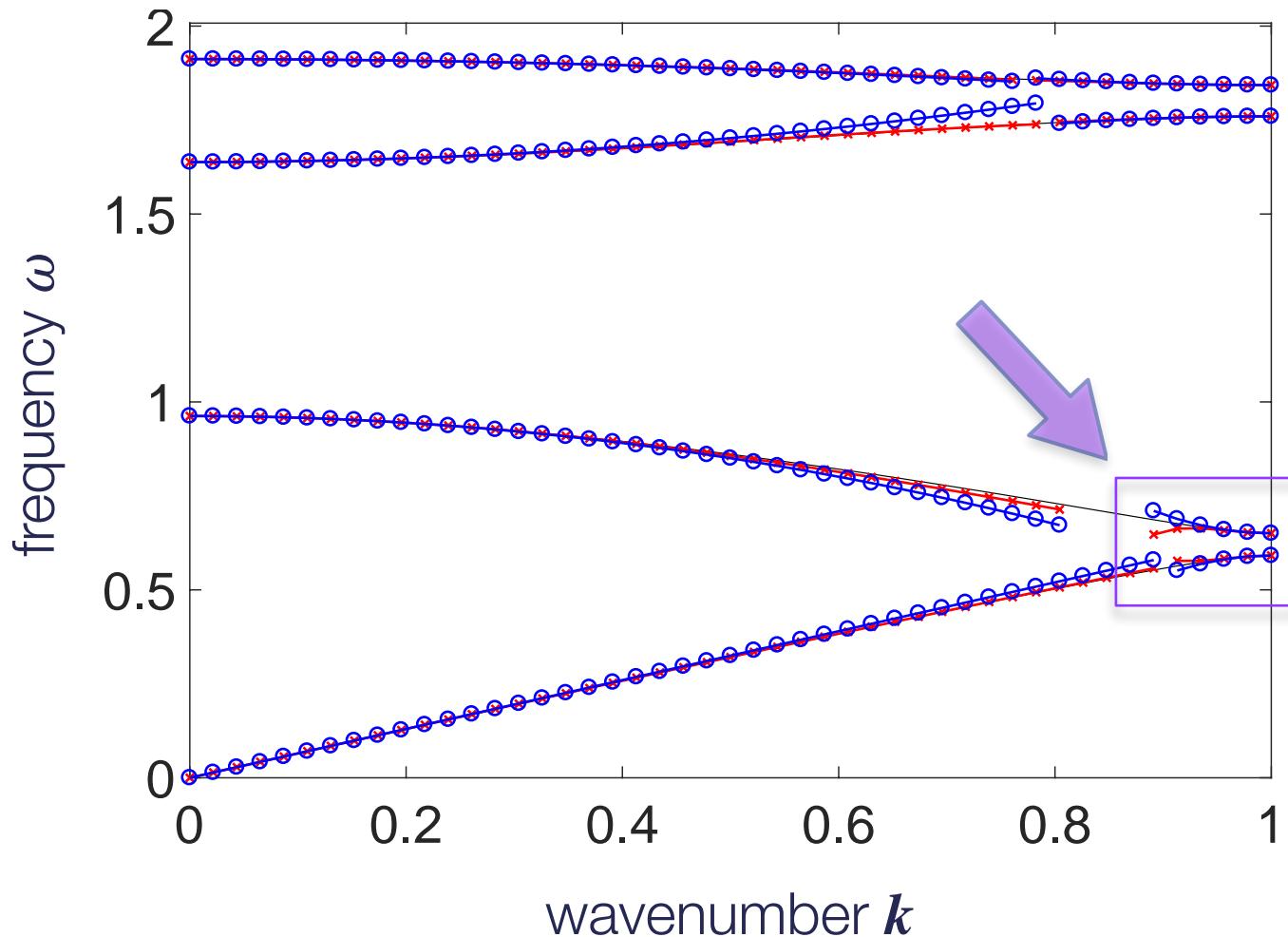
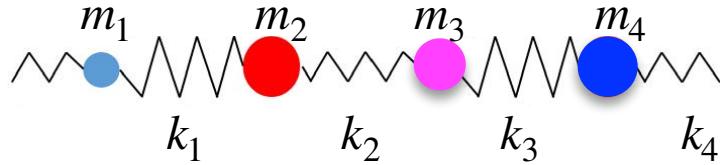
Redux



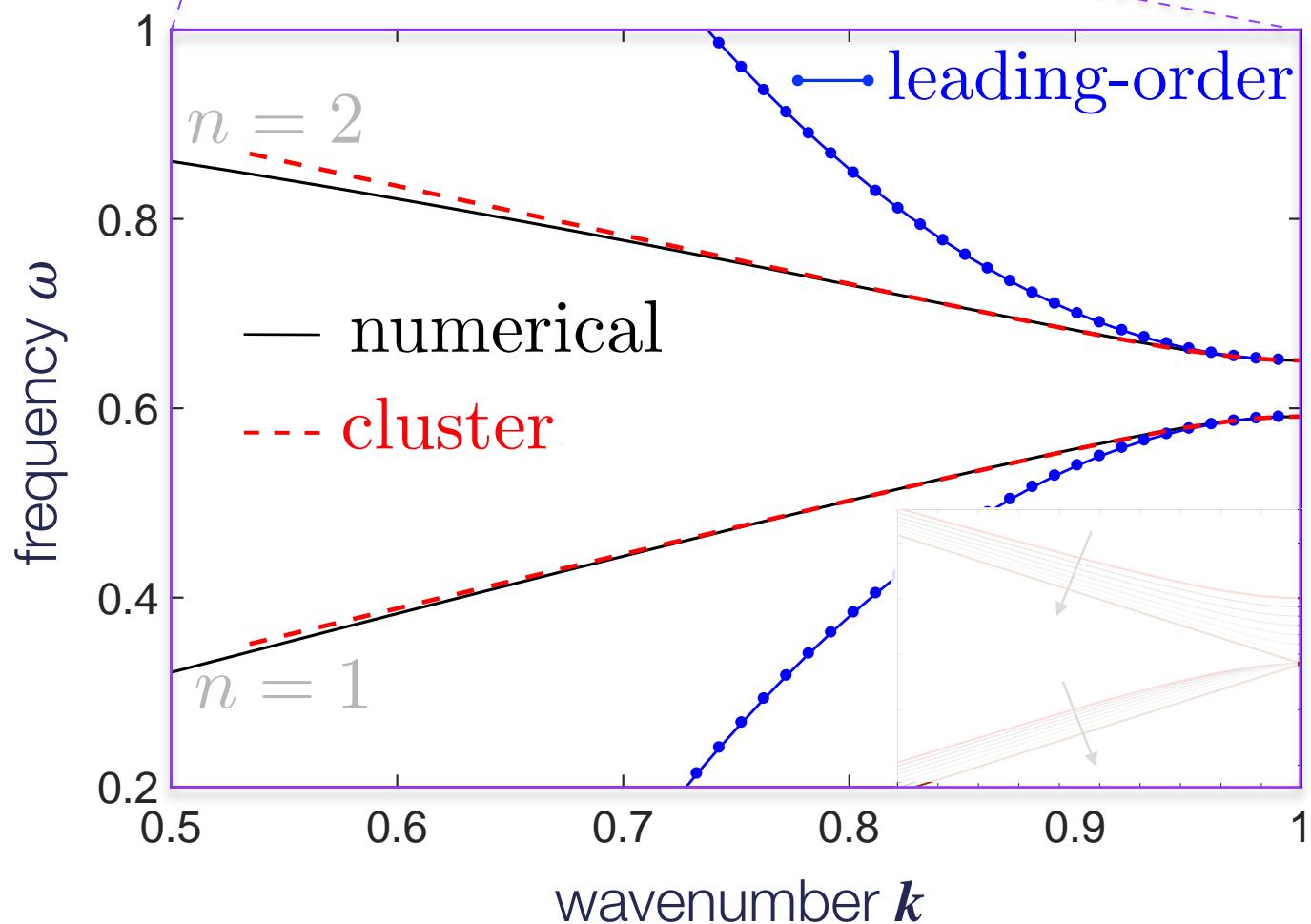
Three nearby eigenvalues



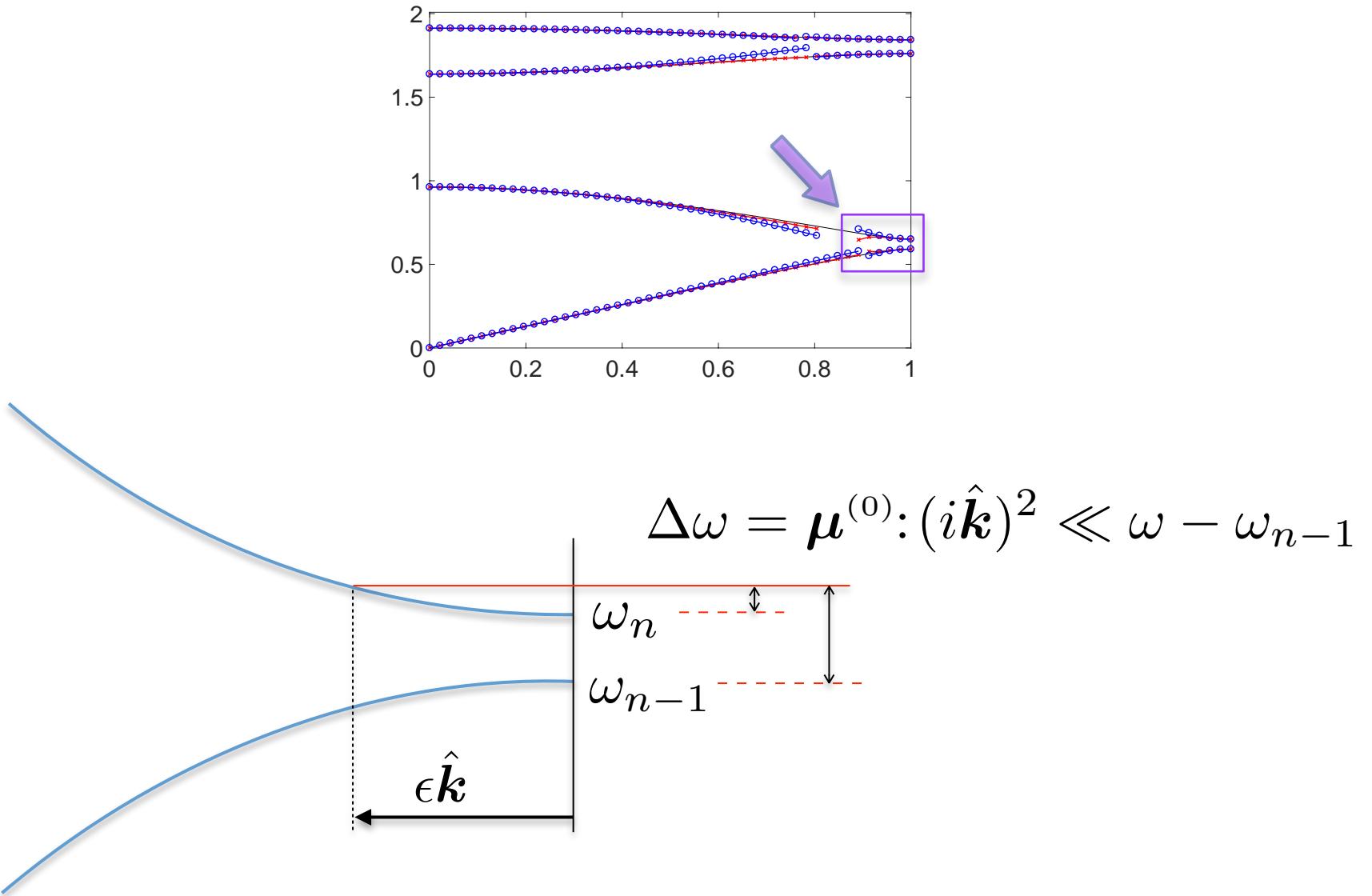
Tetratomic chain



Two nearby eigenvalues

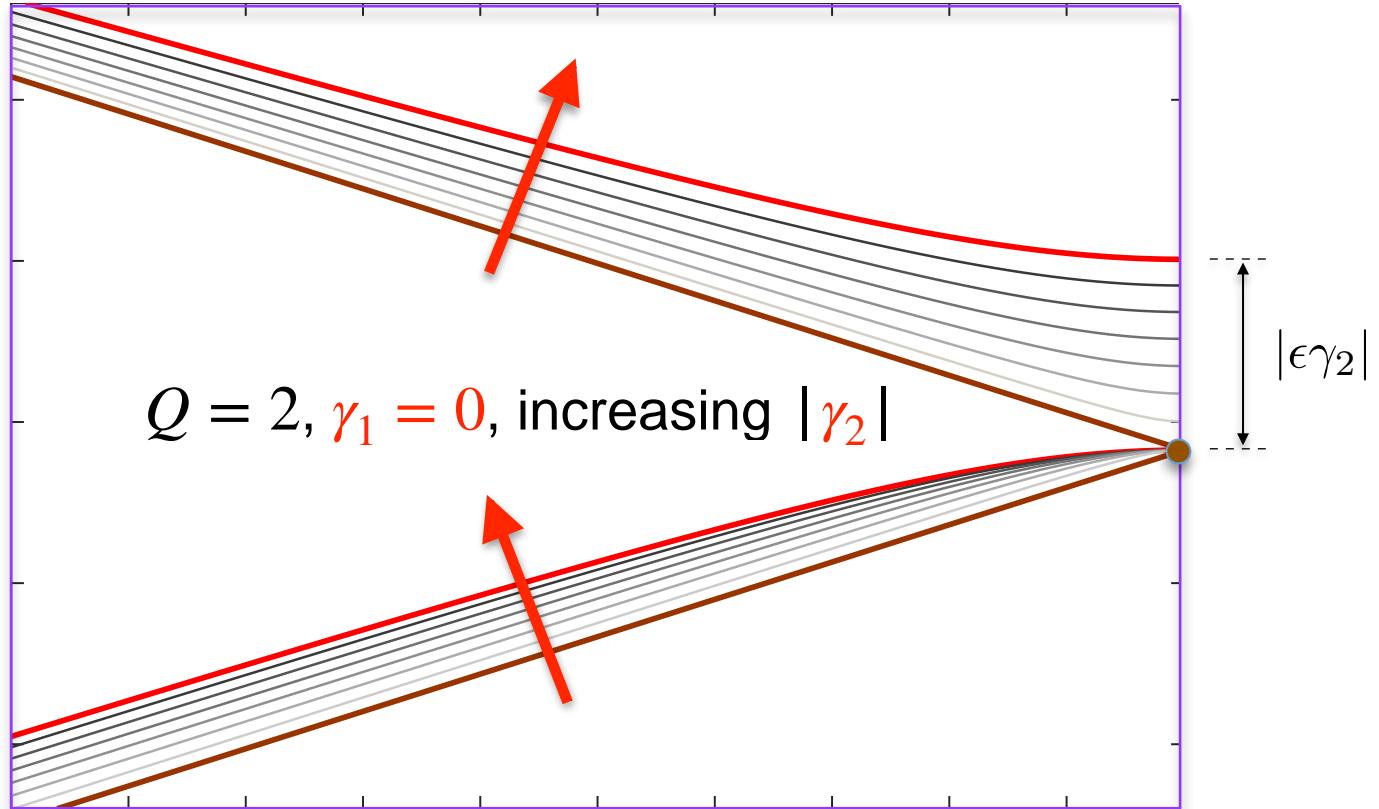


Role of the curvature



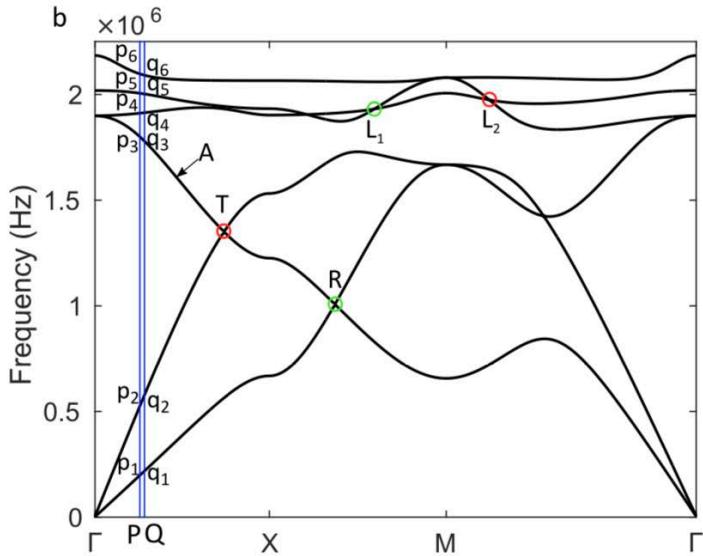
Dirac is (almost) everywhere

$$-\sum_q A_{pq} w_{1q} - \sigma \check{\omega}^2 \sum_q D_{pq} w_{1q} = 0$$



$$\begin{bmatrix} 0 & A_{12} \\ -A_{12} & 0 \end{bmatrix} \xrightarrow{\hspace{1cm}} \begin{bmatrix} 0 & A_{12} \\ -A_{12} & \gamma_2 \rho_2^{(0)} \end{bmatrix}$$

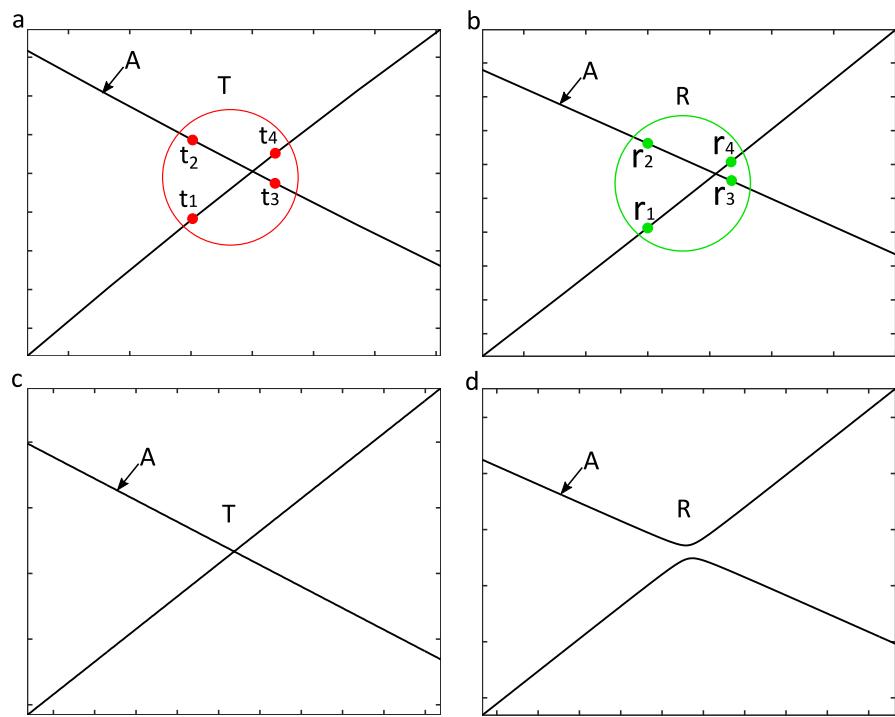
Band sorting & level repulsion



Macroscopic view

?

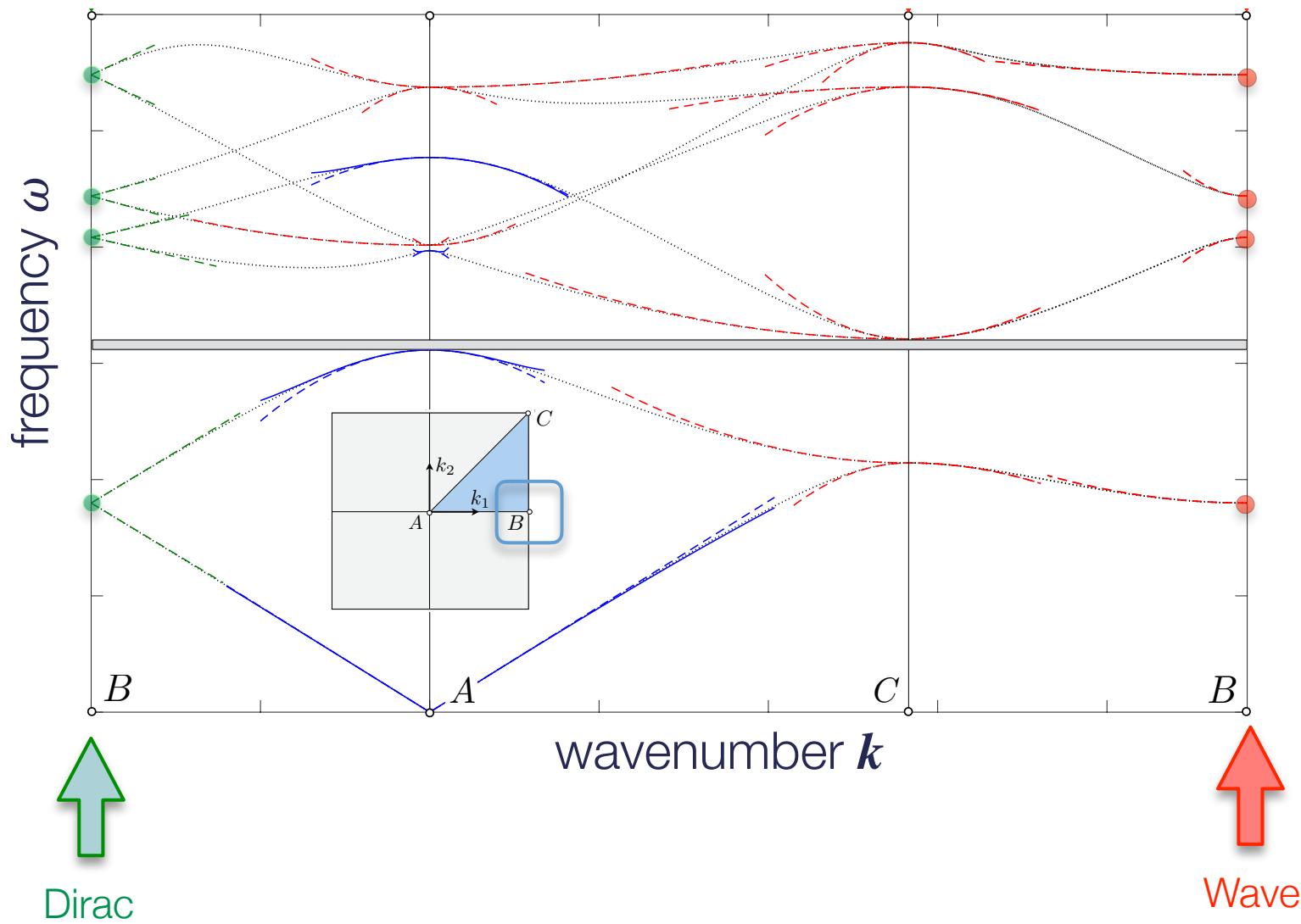
Lu & Srivastava (2018) JMPS, 111



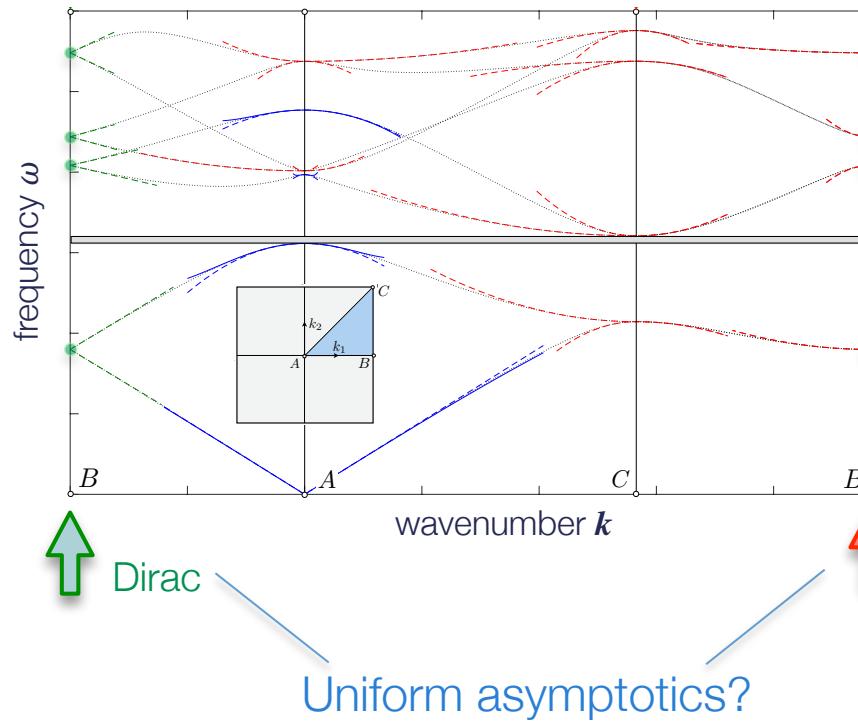
Zoom-in

Remarks

Direction-dependent asymptotic behavior



Remarks



$$\begin{bmatrix} 0 & A_{12} \\ -A_{12} & 0 \end{bmatrix}$$

$$A_{pq} = \theta_{pq}^{(0)} \cdot i\hat{k}$$

$$\theta_{pq}^{(0)} = \langle G\nabla \tilde{\varphi}_{nq}^a \rangle_a^{p\varphi} - \langle G\nabla \tilde{\varphi}_{np}^a \rangle_a^{q\varphi}$$

Lemma 5.11. A sufficient condition for the occurrence of a Dirac point [15], i.e. conical contact between dispersion surfaces, at (k^a, ω_n^a) is given by

$$\text{rank}(A_{pq}) = Q \quad \forall \frac{\hat{k}}{\|\hat{k}\|} \in \{\kappa \in \mathbb{R}^d : \|\kappa\| = 1\}, \quad Q > 1.$$

As a result, 'simple' ($Q = 2$) Dirac points of the scalar wave equation are not possible.

Proof. When $Q = 2$, the eigenvalues of A_{pq} are $\pm \theta_{12}^{(0)} \cdot \hat{k}$, which vanish for $\hat{k}/\|\hat{k}\| \perp \theta_{12}^{(0)}$. ■